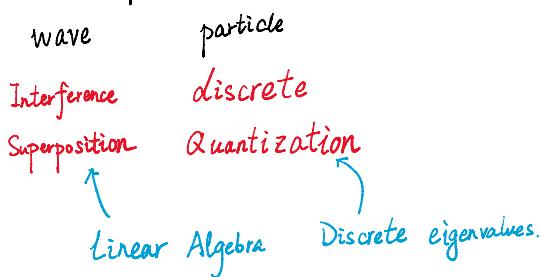


2.1. Formalism of Quantum mechanics.

$\left\{ \begin{array}{l} \text{States to describe systems} \\ \text{evolution} \\ \text{measurement / observables} \end{array} \right.$

1. Basic concept: wave-particle duality



2. States in Quantum mechanics.

Assumption 1: states \Leftrightarrow a vector in Hilbert Space

- vector. $|\psi\rangle$
 $\langle\psi|$ Inner product $\langle\psi|\psi\rangle$
 complex conjugate. norm. $\langle\psi|\psi\rangle \geq 0$
 $|\psi\rangle = \alpha|0\rangle + \beta|1\rangle$

Hermitian:

$$A = A^\dagger$$

$$A^\dagger : \stackrel{\text{(def)}}{=} \langle\psi_1|A|\psi_2\rangle = \langle\psi_2|A^\dagger|\psi_1\rangle^* \quad \left\{ \begin{array}{l} \sum_i |\lambda_i\rangle \langle\lambda_i| = I \\ \langle\lambda_i|\lambda_j\rangle = \delta_{ij} \end{array} \right.$$

Spectral Decomposition

$$\begin{aligned} A &= \sum_i \lambda_i |A_i\rangle \langle A_i| \\ &= \sum_i \lambda_i |\lambda_i\rangle \langle\lambda_i| \quad \text{real} \end{aligned}$$

(see LA 6.4.)

$$|\psi\rangle = \sum_i |\lambda_i\rangle \langle\lambda_i|\psi\rangle$$

3. Observables in Quantum mechanics.

Assumption 2: Observables \Leftrightarrow a Hermitian operator. in Hilbert Space

$$|\psi\rangle \xrightarrow{\text{operator}} A|\psi\rangle$$

operators a chosen basis matrices.

coordinate	X	\hat{X}
momentum	P	$\hat{P} = -i\hbar \nabla$
Energy	$E = E(x, p)$	$\hat{H}(\hat{x}, \hat{p}) = \frac{\hat{p}^2}{2m} + V(\hat{x})$

Hamiltonian operator.

4. Evolution in Quantum mechanics.

$$|\psi(t=t_0)\rangle \rightarrow |\psi(t=t_0+\Delta t)\rangle$$

Assumption 3: Schrödinger Equation.

$$i\hbar \frac{\partial|\psi\rangle}{\partial t} = \hat{H}|\psi\rangle$$

if \hat{H} is time-independent.

$$A = \sum_i \lambda_i |\lambda_i\rangle \langle\lambda_i|$$

$$f(A) = \sum_i f(\lambda_i) |\lambda_i\rangle \langle\lambda_i|$$

$$-i\hat{H}\Delta t \rightarrow -i\hbar\Delta t \cdot \sum_i \lambda_i \langle\lambda_i| e^{-i\frac{\hbar}{\hbar}\Delta t} |\lambda_i\rangle$$

$$i\hbar \frac{d|\psi\rangle}{dt} = \hat{H}|\psi\rangle$$

if \hat{H} is time-independent.

$$|\psi(t)\rangle = e^{-i\hat{H}(t-t_0)/\hbar} |\psi(t_0)\rangle$$

Evolution operator $U(t, t_0)$

$$U^\dagger U = I \quad \leftarrow \begin{array}{l} \text{Unitary} \\ \text{norm-preserving} \end{array} \quad (\text{see LA 6.5}).$$

$$\langle \psi(t) | \psi(t) \rangle = \langle \psi(t_0) | U^\dagger U | \psi(t_0) \rangle = \langle \psi(t_0) | \psi(t_0) \rangle$$

5. Quantum no-cloning theorem

system	A	B
initial	$ \psi\rangle \otimes 0\rangle$	blank state.
COPY	$ \psi\rangle \otimes \psi\rangle$	

Tensor Product.

$$|\psi\rangle \otimes |\psi_2\rangle$$

Proof. If this process is possible

$\Rightarrow \exists$ unitary operator U :

$$\$. \quad U(|\psi\rangle_A \otimes |0\rangle_B) = |\psi\rangle_A \otimes |\psi\rangle_B$$

This COPY should work for any $|\psi\rangle$.

$\Rightarrow |\psi_1\rangle, |\psi_2\rangle$ with $\langle \psi_1 | \psi_2 \rangle \neq 1$ or 0.

$$\Rightarrow \begin{cases} U(|\psi_1\rangle_A |0\rangle_B) = |\psi_1\rangle_A |\psi_1\rangle_B \\ U(|\psi_2\rangle_A |0\rangle_B) = |\psi_2\rangle_A |\psi_2\rangle_B \end{cases} \xrightarrow{\text{(Inner product)}} \underbrace{\langle \psi_1 | \psi_2 \rangle_A \otimes |0\rangle_B}_{=1} = \underbrace{\langle \psi_1 | \psi_2 \rangle_A}_{=0} \underbrace{|0\rangle_B}_{=0}$$

Contradiction!

Great for Security

6. Measurement in Quantum mechanics.

Assumption 4. When measuring A under state $|\psi\rangle$.

the state is projected to an eigenstate $|\lambda_i\rangle$ of A with prob. $P_i = |\langle \lambda_i | \psi \rangle|^2$ and outcome $A = \lambda_i$

Define von Neumann projectors $P_i \stackrel{\text{def}}{=} |\lambda_i\rangle \langle \lambda_i|$

$$\begin{aligned} \text{Initial } |\psi\rangle &\xrightarrow{\quad} \begin{cases} P_i \geq 0 \\ P_i^2 = P_i \\ P_i P_j = \delta_{ij} P_i \\ \sum_i P_i = I \end{cases} \\ \text{Final (After Measurement)} & \end{aligned}$$

$$P_i |\psi\rangle = \langle \lambda_i | \psi \rangle \cdot |\lambda_i\rangle$$